

22

Abstract

23 We investigate the impact of the Pacific Asian Marginal Sea (PAMS) on the second
24 dominant ocean-atmosphere coupled mode variability in the North Pacific through the
25 modification of North Pacific Oscillation (NPO) variability. The second dominant
26 coupled mode reflects the oceanic footprint of Victoria Mode (VM) forced by the
27 meridional variability of NPO. Our observations show that the November meridional
28 wind in the PAMS strongly connects to the southern lobe variability of the matured
29 NPO in January, leading to the following basin scale oceanic VM pattern. No
30 connection was found for the northern lobe of the NPO. Further wave ray tracing
31 analyses confirmed the above atmospheric teleconnection resulting from the Rossby
32 wave propagation in the middle to upper troposphere. The wave trajectories' origin
33 from the PAMS shows a clear eastward propagating pathway to affect the southern
34 lobe of the NPO from the southern lobe of Western Pacific pattern at 500 hPa on the
35 time scale of one month. No ray trajectories from the lower troposphere can propagate
36 eastward to influence the central-eastern subtropical Pacific. Further analysis suggests
37 that the initialization of the East Asian Winter Monsoon may play an important role in
38 controlling low-level meridional wind variability in the PAMS, but does not explain
39 its completed link with the southern lobe of the NPO.

40

41 Keywords: North Pacific Oscillation, East Asian Winter Monsoon, Victoria Mode,
42 Pacific Asian Marginal Seas

43 **1. Introduction**

44 The ocean-atmosphere (O-A) coupled system and its variation in the Pacific show
45 a large impact on global weather and climate. For example, El Niño–Southern
46 Oscillation (ENSO) is a particular mode of Pacific climate variability with strong
47 coupling between the atmosphere and ocean in the tropical Pacific. Several previous
48 studies have found that the teleconnection between the tropics and mid-latitudes may
49 drive ENSO variations through the “Seasonal Footprinting Mechanism” (SFM)
50 proposed by Vimont et al. (2003). The SFM asserts that the second leading pattern of
51 winter atmospheric circulation, the North Pacific Oscillation (NPO), and its variability
52 over the North Pacific can significantly impact spring sea surface temperature (SST)
53 anomalies in the central North Pacific by modifying the wind-stress fields and
54 changing the net surface heat flux over the North Pacific (Vimont et al., 2009;
55 Alexander et al., 2010; Furtado et al., 2012; Ding et al., 2015a; Ding et al., 2015b).
56 Hereafter, the seasons referred to are those occurring in the Northern Hemisphere.
57 The NPO is defined as the second leading mode of winter sea level pressure (SLP)
58 anomalies over the North Pacific (Walker & Bliss, 1932; Rogers, 1981; Vimont et
59 al., 2003; Linkin and Nigam 2008). It consists of a meridional dipole in SLP
60 anomalies, centered near Hawaii and over Alaska. The NPO significantly influences
61 downstream weather and climate conditions over North America during wintertime
62 (e.g., Rogers 1981; Linkin and Nigam 2008). However, its cause and the associated
63 dynamical process modulating its variability are not completely understood.

64 Most previous studies have stated the variability of NPO is the stochastic
65 atmospheric process (e.g., Rogers, 1981; Furtado et al., 2012). Some studies have
66 recently suggested the nonstationary variability of NPO and its possible connection
67 with the East Asian climate. Wang et al. (2007) found that, before 1975, the NPO was
68 mainly influenced by the ENSO teleconnection with anomalous stationary wave
69 propagation from the subtropical-central Pacific to the northern Pacific. However,
70 after 1975, the NPO was associated with a winter circumglobal wave train pattern
71 over the extratropical region in the Northern Hemisphere and was closely linked to
72 the East Asian climate. Pak et al. (2014) further found another nonstationary
73 relationship between the NPO and the East Asian Winter Monsoon (EAWM) with a
74 sudden change after 1988. This can be associated with the pronounced decadal
75 weakening of the Siberian high system over the Eurasian continent and the eastward

76 migration of the NPO dipole. A strong correlation was found during the 1973-87
77 strong EAWM epoch, but vanished during the 1988-2002 weak winter monsoon
78 epoch. These studies suggested the complicated dynamics associated with the origin
79 of NPO and how the NPO can be modulated.

80 Similarly, the impact of the NPO on the ENSO is hardly straightforward. Some
81 studies noted that the existence of negative (positive) NPO in winter(0) does not
82 always result in El Niño (La Niña) the following winter [winter(1)] (Pegion and
83 Alexander, 2013; Ding et al., 2015a). Hereafter, we denote the year in which the El
84 Niño (La Niña) developing year 0 and the preceding and following years as year(-1)
85 and year(1), respectively. The impact of the NPO on the development of ENSO
86 conditions through the SFM may depend on the state of the tropical Pacific (e.g.,
87 Anderson, 2007; Alexander et al., 2010; Ding et al., 2015a). Ding et al. (2015a)
88 further confirmed that there is no direct connection between the NPO and ENSO
89 variations. However, the basin-scale oceanic Victoria Mode (VM, Bond et al., 2003),
90 defined as the second leading mode of north Pacific SST anomalies, may act as a
91 more effective pathway for NPO-like atmospheric variability to drive ENSO
92 variability via the SFM. The well-known Pacific Meridional Mode (PMM, Chiang &
93 Vimont, 2004; Chang et al., 2007), which directly links to the SFM, is merely a partial
94 response of the basin-scale VM (see the detailed discussion in Ding et al., 2015a).

95 Indeed, better connection of the basin-scale VM to the ENSO variation than
96 PMM (or NPO) mode suggests the active role of the Western North Pacific (WNP).
97 Figure 1 shows the correlation map of Niño4 index in January(1) and Pacific SST
98 anomalies at different lags (only correlation coefficients significant at the $p < 0.05$
99 level are shown) from Extended Reconstructed SST version 4 (ERSST.v4)
100 observations (1948–2015). The Niño4 index is chosen here because it closely relates
101 to the central Pacific warming (CPW) commonly excited by the PMM (e.g., Yu and
102 Kim 2011; Kim et al. 2012; Vimont et al., 2014). Similar patterns and correlations can
103 be found using Niño3 or Niño3.4 indices. The evolution pattern shows the different
104 developing stages of VM before the mature phase of ENSO. We note that a good
105 correlation between the SST anomalies and Niño4 index can be identified in the WNP
106 and eastern Indian Ocean more than nine months prior (left panel). The correlation is
107 higher than 0.6 (significant at the $p < 0.05$ level) at the 12-month lead time in the
108 WNP. This high correlation region (signals already emerged at the 15-month lead
109 time) is located at the southern edge of the WNP index defined in Wang et al. (2012)

110 and may directly trigger the winter(0) SST dipole in the WNP, which is related to the
111 development of ENSO in the following winter(1). Similar patterns can also be found
112 in the Hadley Centre SST data set (not shown).

113 The above high correlation in the WNP is supported by several other regional
114 studies that have also noted SST anomalies in the marginal seas of the WNP are
115 colder than normal monthly climatology in the developing years [year(0)] of El Niño.
116 Hong et al. (2001) found that summer(0) SST anomalies in the East (Japan) Sea tend
117 to be colder than during year(-1) prior to the developing year [year(0)]. SST
118 anomalies during El Niño developing years are also opposite to those during La Niña
119 developing years. Similar cold SST anomalies in spring(0) have been further
120 confirmed based on three independent long-term observational stations off the east
121 coast of South Korea (Jo et al., 2014). These signals of cold SST anomalies in the
122 WNP are consistent with the correlation map in Figure 1.

123 In particular, Wang et al. (2012) used the winter(0) SST anomaly dipole to
124 enhance the ENSO forecast. A strong correlation between ENSO and the preceding
125 SST anomalies dipole in the Pacific Asian Marginal Seas (PAMS) of the WNP was
126 found based on a robust statistical analysis (similar to Figure 1). They thought that the
127 spatial pattern in the WNP shares similar characteristics with the MM (Chiang &
128 Vimont, 2004; Chang et al., 2007; Zhang et al., 2009), except that the meridional SST
129 anomalies gradient and low-level zonal wind anomalies occur in the western tropical
130 Pacific. These studies strongly suggested that large-scale SST anomalies over the
131 subtropical/extratropical WNP are significantly related to El Niño 12 months later
132 through the modulation of the winter NPO and can be used as a useful predictor for El
133 Niño development. However, it remains unclear as to how these regional WNP
134 findings are linked to the NPO and the forced basin-scale VM.

135 In this paper, we focus on the pathway modulating the NPO-like atmospheric
136 variability (specifically the southern Lobe) from the PAMS. We isolate the dominant
137 role of the WNP (originating primarily from the PAMS) and its influence on the
138 winter NPO variability. The driving mechanism and pathway modulating the NPO
139 will be addressed by means of observation and the Rossby wave ray tracing analyses.
140 Section 2 introduces the observational data and the Rossby wave ray tracing method
141 theory. Section 3 investigates the modulation of the NPO variability resulting from
142 the WNP in the observation. Section 4 validates the associated pathway using the
143 Rossby wave ray tracing. Section 5 discusses how this PAMS origin explains the

144 ENSO precursor signals in the literature. Finally, conclusions and suggestions for
145 further work are provided in Section 6.

146

147

148 **2. Data, analysis method and theory**

149 **2.1 Observational data**

150 Three observational datasets for the period 1948–2015 are used here. The
151 monthly SLP, surface wind components, surface latent heat flux, geopotential height,
152 and wind vectors are taken from the National Centers for Environmental Prediction –
153 National Center for Atmospheric Research (NCEP–NCAR) reanalysis project (Kalnay
154 et al., 1996; Kistler et al., 2001) on a $2.5^\circ \times 2.5^\circ$ horizontal grid resolution and 12
155 vertical pressure levels ranging from 1000 to 100 hPa. The climatology of 1986–2015
156 is removed. For the surface latent heat flux, we also verify the dynamical processes
157 discussed in this paper using objectively analyzed air-sea fluxes, OAFLUX (Yu et al.,
158 2008), during 1983–2009. Some minor differences can be found in the amplitude;
159 however, there is no significant difference in terms of the patterns comparing with the
160 NCEP–NCAR reanalysis product. Therefore, for the sake of consistency, we present
161 all observational results based on NCEP–NCAR reanalysis throughout the paper. The
162 ERSST.v4 data are taken from the National Climatic Data Center on a $2^\circ \times 2^\circ$
163 horizontal grid (Smith et al., 2008). All relevant SST indices are calculated from the
164 ERSST.v4 data based on the standard definition. The representation of the NPO is
165 defined according to the second dominant SLP anomalies mode in the combined
166 empirical orthogonal function (CEOF) analysis (Tseng et al., 2016a). The linear trend
167 is removed for all datasets.

168 **2.2 Rossby wave ray tracing theory**

169 Rossby wave ray tracing theory can provide comprehensive insights into the
170 wave evolution and has been shown to effectively characterize the subsequent
171 evolution of the Rossby wave train excited by the tropical heating source (Hoskins
172 and Karoly 1981; Shaman and Tziperman 2005; Zhao et al., 2015). Following Li et al.
173 (2015) and Zhao et al. (2015), the dispersion equation for propagation of Rossby
174 wave on a horizontally non-uniform flow with the inclusion of the meridional
175 component may be written as,

176

177
$$\omega = \bar{u}_M k + \bar{v}_M l + \frac{l \partial \bar{q} / \partial x - k \partial \bar{q} / \partial y}{K^2}, \quad (1)$$

178 where ω is the frequency, $(\bar{u}_M, \bar{v}_M) = (\bar{u}, \bar{v}) / \cos \varphi$ is the Mercator projection of the
 179 basic-state zonal and meridional winds, $\bar{q} = 2\Omega \sin \varphi + \nabla^2 \bar{\psi}$ is the basic-state
 180 absolute vorticity, φ is latitude, Ω is the rotation rate of the earth, and $\bar{\psi}$ is the
 181 basic-state streamfunction; $K = \sqrt{k^2 + l^2}$ is the total wavenumber, k and l are the
 182 zonal and meridional wavenumbers, respectively. The wave ray is defined as a
 183 trajectory that is locally tangential to the group velocity. Using the Eq. (1), the group
 184 velocity of the Rossby wave packet takes the form

185
$$u_g = \bar{u}_M + \left[(k^2 - l^2) \partial \bar{q} / \partial y - 2kl \partial \bar{q} / \partial x \right] / K^4 \quad \text{and} \quad (2a)$$

186
$$v_g = \bar{v}_M + \left[2kl \partial \bar{q} / \partial y + (k^2 - l^2) \partial \bar{q} / \partial x \right] / K^4. \quad (2b)$$

187 The evolutions of wavenumbers k and l along a trajectory are determined by
 188 kinematic wave theory (Whitham 1960), as follows

189
$$\frac{d_g k}{dt} = -k \frac{\partial \bar{u}_M}{\partial x} - l \frac{\partial \bar{v}_M}{\partial x} + \frac{k \partial^2 \bar{q} / \partial y \partial x - l \partial^2 \bar{q} / \partial x^2}{K^2} \quad \text{and} \quad (3a)$$

190
$$\frac{d_g l}{dt} = -k \frac{\partial \bar{u}_M}{\partial y} - l \frac{\partial \bar{v}_M}{\partial y} + \frac{k \partial^2 \bar{q} / \partial y^2 - l \partial^2 \bar{q} / \partial x \partial y}{K^2}, \quad (3b)$$

191 where $d_g / dt = \partial / \partial t + u_g \partial / \partial x + v_g \partial / \partial y$ represents the material derivative moving
 192 with the group velocity.

193 For each initial zonal wave number k and an initial position within the region
 194 where the wave source has been perturbed, we use Eq. (1) to solve for the initial
 195 meridional wavenumber l by assuming $\omega = 0$ (stationary waves). Then Eqs. (2) and
 196 (3) were integrated to derive the ray trajectory; they were terminated when the spatial
 197 scale of Rossby waves was less than 1000 km. The background climatological mean
 198 wind flow is based on the November-December-January (NDJ) NCEP-NCAR
 199 reanalysis for the period of 1963-2015. Since the Wentzel-Kramers-Brillouin
 200 approximation used in the Rossby wave theory also requires a slow-varying
 201 background flow (Li et al., 2015), we also smoothed the background flows of \bar{u} and
 202 \bar{v} by multiplying a spherical harmonic component with total wavenumber n by
 203 $\exp\{-\xi [n(n+1)]^2\}$, where the coefficient ξ has been set such that the harmonic
 204 component of $n = 16$ was reduced by 37%. This reduced the potential impacts of

205 the small-scale structures in the basic state. Further examination indicates that the
206 results are not sensitive to the smoothing method applied. Further details can be found
207 in Li et al. (2015) and Zhao et al. (2015).

208

209 **3. Link between the surface meridional winds in the PAMS and the NPO**

210 Figure 1 suggests the ENSO variability is connected with the precursory spring
211 Pacific VM and its atmospheric forcing of NPO (Ding et al., 2015a). These surface
212 expressions can be well explained using the second dominant CEOF (CEOF2) of SLP
213 and SST anomalies in the Pacific (Figure 2). The meridional dipole structure of SLP
214 anomalies in the North Pacific is the typical signature of the NPO pattern (Figure 2a).
215 This NPO pattern is very similar to that reported earlier (Linkin & Nigam, 2008;
216 Furtado et al., 2012) and is a robust winter atmospheric feature. There is another weak
217 pressure low in the subtropical WNP, which has drawn only minor attention
218 (Anderson, 2007), but may be important to modulate the NPO variability (discussed
219 in section 5). The CEOF2 of the SST anomalies resembles the VM described above
220 (Bond et al., 2003; Di Lorenzo et al., 2008; Ding et al., 2015a)—a region of negative
221 SST anomalies extending from the WNP to the Kuroshio Extension, encircled by
222 warm SST anomalies around the North Pacific coast reaching the central tropical
223 Pacific (Figure 2b). The variability of CEOF2 is directly connected to the most recent
224 warm Blob observed in the Northeastern Pacific evolved from late 2013 to 2016 and
225 the persistence of the 2014/15 North Pacific climate linking to the strong 2015 El Niño
226 (Di Lorenzo and Mantua, 2016; Tseng et al., 2016a). However, these recent studies
227 only address the evolution of the VM in the El Niño events and the multi-year
228 enhancement (Figure 1c-1f). They did not address the preceding SST anomalies
229 observed in the PAMS [winter(0)] one year in advance of the occurrence of El Niño
230 (Figure 1b), which may link to the development of the NPO and the following
231 evolution of VM.

232 Based on the observational evidences (e.g., Wang et al. [2012] and Figure 1), we
233 propose the local atmospheric heating related to the SST dipole in the PAMS can
234 affect the NPO developing phase through the atmospheric teleconnection during the
235 early winter (i.e., November-December-January, NDJ). The teleconnection can be
236 initiated from the change of alongshore wind in the East Asian (discussed in section
237 5), particularly the meridional wind component. Figure 3 shows the correlation map
238 of the November surface meridional wind anomaly in the PAMS (120~140°E,

239 5~25°N), hereafter referred to as “ V_{PAMS} ”, with several lags of SLP anomalies (0 to 3
240 months). Prior to the correlation analysis, the Niño3.4 index was linearly removed
241 from the wind fields using the regression. The removal of the Niño3.4 index from the
242 data ensures that the surface meridional winds are independent of any ENSO
243 variability that occurs simultaneously in the tropical Pacific. This region is chosen
244 from the largest SST anomalies correlation map 12 months prior to the mature phase
245 of ENSO shown in Figure 1b. The results are not sensitive if a slightly larger region is
246 selected (e.g., 100~140°E, 5~30°N) as long as the high correlation area in Figure 1b is
247 included. The simultaneous large correlation can be found along the coast of China
248 from South China Sea (SCS) to Yellow Sea (maximum negative correlation $R=-0.73$),
249 indicating the strong geostrophic wind balance. Another comparable largest
250 correlation can be found when the November V_{PAMS} leads the SLP anomalies by 2
251 month ($R=0.56$) in the region just above the southern branch of the VM (defined as
252 SLP_s , 180°~140°W, 5~25°N, see the grey box in the Figure 1c). This is the region
253 where the southern lobe of the NPO starts to initiate the VM or MM through the SFM
254 (Ding et al., 2015a). Furtado et al. (2012) suggested a large distinction in the
255 dynamics between the northern and southern lobes of the NPO while the southern
256 lobe of the NPO likely linked to the CPW. This viewpoint is supported by the
257 significant correlation in Figure 3c, while the correlation is not significant for the
258 northern lobe of the NPO. Our results suggest a more direct impact of the PAMS
259 surface wind variability in modulating the NPO within two months than the
260 low-frequency influence proposed by Furtado et al. (2012).

261 The two-month lag near the southern lobe of the NPO is not a coincidence.
262 Figure 4a verifies the observed lead-lag correlation between the V_{PAMS} and the SLP
263 anomalies in the southern branch of the VM region (the grey box in the Figure 1c),
264 i.e., influencing the southern lobe of the NPO. Our results suggest that any potential
265 heating source in the PAMS can affect the central-subtropical Pacific two months later.
266 Another significant, but weaker, correlation can be found in March when SLP
267 anomalies in the southern branch of the VM region has no lag with the V_{PAMS} . This
268 simultaneous correlation is actually associated with the spring footprint of VM
269 resulting from the NPO forcing and supports the main difference between the VM and
270 MM (discussed in Ding et al., 2015a). Figure 4b shows a similar lead-lag correlation
271 between the V_{PAMS} and the PC2 associated with CEOF2 (NPO/VM) pattern described
272 above. Note that a minus sign is applied to the PC2 to ensure the consistency of the

273 correlation sign. As expected, a clearer correlation can be found, indicating the PC2
274 actually leads the V_{PAMS} in spring, causing the basin-scale VM pattern. The triggering
275 of NPO on the ENSO development is not only through the MM but also affects the
276 surface wind in the PAMS. Since this spring-to-summer evolution has been detailed
277 previously, our current study emphasizes the winter development of the NPO.

278 The correlation analysis suggests the enhanced November V_{PAMS} in the PAMS
279 leads to the growth of winter NPO, reaching its maximum in January (particularly the
280 southern lobe). The NPO then forces the surface SST footprint of VM through the
281 SFM, which matures in the following March or April (Ding et al., 2015a). Figure 5a
282 shows the time series of the November $-V_{\text{PAMS}}$, January $-\text{SLP}_s$ and the March PC2,
283 respectively. Negative signs of V_{PAMS} and SLP_s are included here to be consistent
284 with the phase of CEOF2 and its associated PC2 shown in Figure 2. The variation
285 between the November $-V_{\text{PAMS}}$ and January $-\text{SLP}_s$ matches reasonably well,
286 especially after 2000. The corresponding years which have magnitudes greater than
287 one standard deviation are also listed in Table 1.

288 Figure 6 further shows the correlation map between the November V_{PAMS} with
289 the geopotential height anomalies (top three panels: 200 hPa, 500 hPa, and 1000 hPa,
290 respectively) for November and two consecutive lag months (December and January).
291 The simultaneous large positive/negative correlation in November is clear at 1000 hPa
292 along the coast of China, covering Southeast Asia (consistent with Figure 3), while
293 the positive correlation offshore in the WNP can further extend up to 200 hPa. This
294 suggests a northward wind in the PAMS can be associated with a zonal dipole of
295 cyclonic and anticyclonic eddies near the surface, resulting from a local heating
296 source (the same correlation map with latent heat flux is also shown at the bottom).
297 The upward extension to the mid and upper troposphere can be supported by the
298 similar correlation map with the vertical velocity, Ω , in November (upward vertical
299 motion corresponds to the negative Ω) in Figure 7. The upward motion is also
300 consistent with the region of negative OLR anomalies above the PAMS in Figure 8a,
301 corresponding to the anomalously strong convection and rainfall. The phenomenon is
302 reversed for the equatorward meridional velocity in the PAMS (i.e., $-V_{\text{PAMS}}$). These
303 results suggest a barotropic structure is formed in November associated with the
304 surface meridional wind in the PAMS, which is connected to the January NPO pattern
305 in the eastern North Pacific.

306 This atmospheric teleconnection may likely result from the Rossby wave train. In

307 December, the remote atmospheric teleconnection associated with the November
308 V_{PAMS} starts to emerge and evolve in the eastern North Pacific (middle panel in Figure
309 6). The largest correlation occurs in the correlation map of January over the southern
310 branch of the VM in the lower troposphere. This is also associated with a positive
311 correlation in the Ω at 500 hPa and 200 hPa (right panel in Figure 7), suggesting a
312 strong response on the SLP anomalies. These results show a propagating atmospheric
313 teleconnection in the mid to upper troposphere which affects the southern lobe of
314 NPO.

315

316 **4. Rossby wave train mechanism**

317 In order to further explore the described pathway to modulate the winter NPO
318 from the PAMS, we used a theoretical wave ray tracing approach to analyze the
319 potential atmospheric teleconnection resulting from the PAMS during the NDJ. These
320 theoretical analyses examine if the subsequent evolution of the Rossby wave train
321 excited by the heating source from the PAMS can further influence the eastern North
322 Pacific (thus the strength of NPO). Figure 9 shows the wave ray trajectories initiated
323 by the heating sources (crosses) over the PAMS regions under the climatological 500
324 hPa basic flows and integrated for 20 days. The initial position arrays ($110^{\circ}\text{E}\sim 140^{\circ}\text{E}$
325 and $5^{\circ}\text{N}\sim 30^{\circ}\text{N}$, slightly larger than the domain specified in Figure 1b) over the PAMS
326 region were selected to represent the impact of the source locations on the wave
327 propagation. The results shown here are for both positive and negative initial
328 meridional wavenumber l to reflect initially northward and southward propagation
329 of stationary waves, respectively.

330 Interestingly, the sources over the northern PAMS region (north of 20°N) can
331 indeed excite two branches of waves that propagate eastward to the North Pacific
332 sector: the northern branch is poleward and reflected equatorward at mid-latitudes
333 (about 60°N for $k=2$). The other southern branch is equatorward and reflected
334 poleward at the latitudes where the zonal winds equal zero ($U = 0$; about 20°N). A
335 major part of the northern branch finally arrives in the central-eastern North Pacific
336 over the southern branch of the VM before moving eastward to the American
337 continent. Trajectories based on the initial zonal wavenumbers $k=4$ and $k=6$ exhibit
338 qualitatively similar paths, except waves with a larger zonal wavenumber can
339 propagate higher latitudes before equatorward reflection. These two branches of

340 waves confine the wave activity to the North Pacific sector and confirm the sources
341 over the PAMS region can potentially modulate the NPO variability after 20 days,
342 specifically the southern lobe over the southern branch of the VM. This
343 central-eastern North Pacific region is also consistent with the largest correlation
344 between November V_{PAMS} and the January (also December but weaker) geopotential
345 height at 500 hPa shown in Figure 6. Based on the NDJ climatological winds, our
346 integration of wave ray equations indicates the energy transfer from the PAMS source
347 locations can reach the south lobe of NPO in 20 days, which approximately resemble
348 to the 1-2 month lag shown above.

349 In contrast, the sources over the southern part of the PAMS region (south of 20°N)
350 cannot excite waves that can spread eastward. All trajectories are moving westward
351 into tropical western Pacific. This suggests that the north-south locations in the PAMS
352 are critical to determine their impacts. Only the sources within the southern lobe of
353 Western Pacific (WP) teleconnection pattern (i.e., north of about 20°N) (Wallace and
354 Gutzler, 1981; Chow et al., 2016) can influence the Southern lobe of NPO, consistent
355 with the strong connection between the WP and NPO patterns (Linkin and Nigam,
356 2008).

357 Similar wave ray trajectories can also be found at 250 hPa (Figure 10) but the
358 number of trajectories reaching the central-eastern North Pacific is reduced. More
359 trajectories recirculate anti-cyclonically at 250 hPa than 500 hPa for $k=2$ (i.e.,
360 belonging to the sources over the southern part of the PAMS). This is mainly because
361 the vertical structure of the WP tilts northward with increased altitudes (similar to the
362 November correlation map shown in the left panel of Figure 6). Also, the wave
363 trajectories affecting the southern lobe of the NPO are reflected at lower latitudes at
364 250 hPa than 500 hPa, as the jet stream is stronger in 250 hPa leading to the lower
365 turning latitudes.

366 The trajectories are entirely different in the low troposphere. Nearly all rays
367 propagate westward and cannot propagate to the central-eastern North Pacific under
368 the 850 hPa circulation condition (Figure 11). This results from the enhanced
369 southwestward EAWM during the NDJ month. In general, our wave ray tracing
370 analysis, consistent with the NDJ observation shown in section 3, provides evidences
371 that the source in PAMS can significantly contribute to the variability of NPO pattern
372 through the Rossby wave propagation. We find the impact is strongest between

373 600-400 hPa during the NDJ background flow configurations and is directly related to
374 the southern lobe of winter WP pattern. We note that only the stationary waves are
375 discussed here because of the climatological winds, where the disturbances propagate
376 as the group velocity with zero phase velocity stationary.

377

378 **5. Discussion**

379 A new pathway was proposed and verified between the PAMS and the NPO
380 during the winter, one that offers a new perspective on the assumed stochastic
381 dynamics of NPO. Past studies on the NPO have focused mainly on the impacts of its
382 teleconnection patterns on the weather and climate of North America (Linkin &
383 Nigam, 2008) and how the NPO initiates ENSO through the SFM (Ding et al., 2015a;
384 Ding et al., 2015b). Most studies looking at the NPO-induced SFM focus on the
385 consequent changes of surface and subsurface heat content in the central-eastern
386 tropical Pacific (e.g., Anderson, 2004; Anderson & Maloney, 2006; Anderson, 2007).
387 Only a few studies have emphasized the origin of NPO variability. The modulation of
388 NPO due to the V_{PAMS} we describe here mostly affects the NPO's southern lobe (near
389 Hawaii over the southern branch of the VM) prior to the NPO-induced SFM. This is
390 confirmed by the observation and the theoretical wave ray tracing analyses. The
391 influence is similar to the finding of low-frequency forcing of CPW in modulating the
392 low-frequency mode of the NPO in Di Lorenzo et al. (2010) and Furtado et al. (2012).
393 However, we point out a more direct impact of the PAMS surface wind variability in
394 modulating the southern lobe of the NPO through the Rossby wave train propagation
395 on the time scale of one month.

396 Our results confirmed that the variability of January NPO southern lobe can be
397 controlled by the November V_{PAMS} in the PAMS. This regional surface wind change
398 is highly associated with the SST anomalies heating source in the PAMS ($R=0.58$).
399 Similar lag correlation maps based on the November SST index defined in Wang et al.
400 (2012) are superimposed in Figure 3 as thick dashed contours (blue and red dashes
401 represent $R=0.5$ and 0.4 , respectively). The spatial patterns are quite similar to the
402 maps based on the November V_{PAMS} . However, the impact of SST anomalies on the
403 NPO southern lobe appears weaker than the impact of November V_{PAMS} (lower
404 correlation). Wang et al. (2012) implied that the surface winds in the WNP occur
405 almost at the same time as the NPO from December to February (Figure 4 in Wang et
406 al., 2012). Based on the monthly observational data (instead of three-month mean) in

407 section 3 and wave ray tracking analyses in section 4, we further clarify that the NPO
408 can be intensified through the middle to upper tropospheric Rossby wave propagation
409 from the WNP. The November V_{PAMS} plays an important role in enhancing the growth
410 of the southern lobe of the NPO, thus forcing the following basin-scale SST footprint
411 of spring VM in the North Pacific. This is indeed consistent with the tight connection
412 between the WP and NPO (Linkin and Nigam 2008). Recent analysis suggested the
413 WP and NPO to be essentially the same variability, with the NPO being the sea level
414 pressure reflection of the WP geopotential height pattern (Nigam 2003; Linkin and
415 Nigam 2008; Nigam and Baxter 2015). Here, we further demonstrate the origin of the
416 NPO and WP variability is linked with the heating source in the PAMS.

417 The V_{PAMS} in the PAMS can be viewed as the trigger-forcing of the tropical–
418 extratropical teleconnection associated with the NPO variability through the Rossby
419 wave propagation. However, many factors can influence its variability. Further
420 analysis of the 11-yr running correlation coefficient between the V_{PAMS} and southern
421 lobe of NPO variability confirmed that the two indices were well correlated after late
422 1970s with increasing correlation after 2000 (Figure 5b) and poorly correlated during
423 early 1960s and mid-1970s, consistent with the nonstationary connection with the
424 EAWM and the intensified circumglobal wave train activity after 1975 (Wang et al.,
425 2007; Pak et al., 2014). The identified PAMS is a region where the thermal contrast
426 between the East Asian continent and the Pacific is important while the EAWM plays
427 a pivot role in the winter climate.

428 We found the response of November V_{PAMS} that we define here varies
429 simultaneously with the SCS winds and extends northward to the east edge of the
430 Ryukyu Islands with a correlation larger than 0.85 (Figure 8b). A one-month lag
431 correlation map of December suggests the November V_{PAMS} intensifies the local
432 cyclonic circulation over the maritime continent (Figure 8c), thus leading to the
433 eastern Pacific teleconnection pattern in January (Figure 8d). It also lags behind the
434 East China Sea and the Japan Sea winds along the coast of China (25~50°N) by 1
435 month ($R>0.5$). These results suggest that the along-coast wind (mostly meridional
436 wind) associated with the initiation of EAWM in the PAMS during the early winter
437 may modulate the NPO variability. We could not find any available EAWM indices
438 well correlated with the V_{PAMS} . The matured EAWM may have a different inherent
439 dynamics so that the EAWM may not be the only dynamic process to control the
440 V_{PAMS} .

441 In addition, the formation of V_{PAMS} (or the along-coast wind) may be partially
442 caused by the shape of the East Asian coasts and partially related to the pressure
443 gradient between the Siberian high and the Maritime Continent low (Wang & Chen,
444 2014). The variability of subtropical high in the WNP may also contribute to V_{PAMS} . It
445 is possibly the variability of V_{PAMS} is the result of the competition between the
446 EAWM and the WNP subtropical high. Unfortunately, we have not yet found a
447 well-defined WNP subtropical high index to confirm this. The circulation variability
448 near the WNP subtropical high may also connect with the triggering pressure low
449 shown in the CEOF2 pattern (Figure 2a) (also found in Anderson, 2004; 2007). The
450 simultaneous November correlation map of V_{PAMS} also shows such a high correlation
451 feature collocated with the pressure low anomaly (Figure 6g). In terms of the link to
452 the following winter(1) El Niño development, it is likely the ocean dynamics,
453 combined with the impact of initiating EAWM, may change the local O-A interaction,
454 thus contributing to the enhanced $-V_{\text{PAMS}}$ through the latent heat release (Figure 6j).
455 For example, the preconditioned Warm Water Volume (WWV) in the recharge–
456 oscillator mechanism reflects the associated oceanic responses and the subsurface
457 changes to precondition the Pacific WWV (as do WWB, e.g., Fedorov, 2002) more
458 than one year ahead of the ENSO events (Chen et al., 2015).

459 The development of ENSO events may require both the trigger from the south
460 lobe of the NPO and the ocean subsurface preconditioning (Tseng et al., 2016b). In
461 general, the local O-A interaction in the PAMS due to the underlying ocean
462 precondition resulting from the Pacific Subtropical Cell circulation may be a key
463 triggering mechanism that affects V_{PAMS} , which then modulates the variability of
464 southern lobe of NPO. But this process requires a more detailed investigation of the
465 subsurface dynamics.

466 There are some additional mechanisms that could potentially excite the
467 thermodynamic heating/cooling source in the PAMS to affect ENSO variability. For
468 example, Chang et al. (2009) found that significant negative (positive) SST anomalies
469 in the WNP appear in the strong (weak) Subtropical Mode Water (STMW) case. Their
470 statistical analysis indicated that summer(-1) STMW variability can also affect ENSO
471 events 18 months later. They suspected that summer STMW(-1) variability produces
472 subtropical atmospheric variability through long-term persistent SST anomalies over
473 its reemergence area to derive the VM pattern, and eventually modulates the
474 amplitude of ENSO events; however, the role of the STMW still needs to be clarified.

475 Moreover, our results are consistent with previous evidence linking the Tsushima
476 Warm Current to the NPO reported by Hirose et al. (2009). They showed that changes
477 in the Tsushima Warm Current (a current fed by the Kuroshio, which passes between
478 Korea and Japan) during autumn were correlated with the following winter NPO/WP
479 pattern. The surface wind stress in the PAMS indeed controls the transport of
480 Kuroshio and its branches due to the Sverdrup circulation. The consequent response
481 of atmospheric circulation to the Kuroshio can also be found through a positive
482 feedback process in the WNP (Qiu & Chen, 2010; Shen et al., 2014; Chow et al.,
483 2016). The intensive coupled O-A interaction in the WNP significantly modulates the
484 North Pacific climate variability.

485

486

487 **6. Conclusion and future work**

488 The origin of the well-known extratropical NPO forced upon ENSO is
489 investigated and its possible pathway originating from the PAMS have been analyzed
490 in this study using the observation and the wave ray tracing analyses. The NPO
491 commonly forces the basin-scale VM, leading to a matured O-A coupled mode of
492 variability in the subtropical North Pacific (i.e., CEOF2). The proposed pathway well
493 explains all northern hemispheric precursors of El Niño in the literature related to the
494 NPO/VM and the link of PAMS one year in advance (Figure 1). We confirm the
495 origin from the V_{PAMS} modulating the NPO/VM variability. The schematic in Figure
496 12 summarizes a possible explanation of dynamical processes linking the PAMS
497 origin and the onset of El Niño. The variability of the southern lobe of the NPO is
498 initiated by the change of V_{PAMS} associated with the surface heating, and then the
499 reduced deep convection propagates northeastward through the Rossby wave
500 propagation, thus enhancing the deep convection (also reducing the subsidence) near
501 the southern branch of the VM. This is followed by a typical negative phase of NPO
502 forced VM in the Northern Hemisphere (Ding et al., 2015a; Ding et al., 2015b)
503 forming the combined SLP and SST anomalies pattern of CEOF2 from spring(0) to
504 summer(0). The enhanced westerly wind can further lead to the onset of El Niño by
505 the changing of Walker circulation.

506 We found the enhanced (reduced) November deep convection associated with
507 V_{PAMS} aligns well with the changes of southern lobe of the November WP and
508 January subsidence in the central-eastern subtropical Pacific. These processes are

509 robust regardless of the length of the analysis years. This directly supports the
510 integrated picture of the WP and NPO characterized in Linkin and Nigam (2008) but
511 we specifically point out their WNP origin. The wave propagation is on a one-month
512 time scale. We found no evidence of the impact of NPO on the WNP until spring
513 when the basin-scale VM emerges (Ding et al., 2015a).

514 The wave ray tracing analyses confirm the proposed pathway described in the
515 observation. The wave rays origin from the PAMS show clear eastward propagating
516 trajectories reaching the central-eastern subtropical Pacific over the southern branch
517 of the VM (near the southern lobe of the NPO) in the middle to upper troposphere,
518 similar to the typical Rossby wave propagation. These are totally opposite to the wave
519 ray trajectories in the lower troposphere, which flow mostly westward and never
520 reach eastern Pacific. Both the observation and wave ray tracing analyses show no
521 link of the PAMS heating source with the northern lobe of the NPO in the dynamical
522 processes while the abovementioned atmospheric teleconnection affects primarily on
523 the southern lobe of the NPO.

524 The initiation of EAWM seems to play an important role in controlling a part of
525 the variability of V_{PAMS} but does not explain its completed link with the southern lobe
526 of the NPO. The nonstationary relationship between the EAWM and NPO can also be
527 connected through the variability of V_{PAMS} . Our results suggest a seasonally
528 dependent O-A interaction in the PAMS, possibly linked to the change of subtropical
529 high in the WNP. The November V_{PAMS} acts as a pivotal driver to modulate the
530 NPO/VM (meridional variability) through the atmospheric teleconnection. However,
531 the detailed O-A interaction involved is still not clear, specifically the heating source
532 of SST anomalies. Further investigations into the relevant O-A interaction in the
533 PAMS are planned using the linear baroclinic models and CESM-forcing sensitivity
534 studies.

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Table 1: corresponding years of November - V_{PAMS} , January - SLP_s and March PC2 which have magnitudes greater than one standard deviation

	>1 standard deviation	<-1 standard deviation
November - V_{PAMS}	1958, 1961, 1968,1975, 1978, 1988, 1989, 1992, 1994, 2004, 2009	1950, 1952, 1953,1954, 1966, 1986, 1987, 1997, 1998, 2006, 2008, 2011
January - SLP_s	1956, 1959, 1968, 1969, 1971, 1972, 1980, 1990, 2005, 2010	1951, 1953, 1955, 1958, 1964, 1984, 1987, 1999, 2001, 2009, 2012, 2016
March PC2	1956, 1957, 1963, 1965, 1968, 1969, 1981, 1985, 1991, 1997, 2005, 2014, 2015	1953, 1983, 1987, 1998, 1999, 2000, 2007, 2008, 2016

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669

670 Figure captions

671 Figure 1. The correlation of Pacific SST anomalies (based on 3-month running mean)
672 and January(1) Niño4 index at different lags from 1948–2015. The Niño4 index has
673 lags of (a) 15 months; (b) 12 months; (c) 9 months; (d) 6 months; (e) 3 months; and (f)
674 0 months (no lag). Only $p < 0.05$ is shown. The contour lines correspond to $R=0.6$ (bold
675 black), 0.5 (thin black) and 0.4 (grey). The grey box in (b) shows the specific PAMS
676 region and the grey box in (c) shows the southern branch of the VM.

677 Figure 2 The CEOF2 of SLP (hPa) and SST ($^{\circ}\text{C}$) anomalies and the corresponding PC2
678 in the North and tropical Pacific Ocean. (a) SLP; (b) SST; (c) PC2.

679 Figure 3. Correlation map of the November surface meridional wind anomaly (V_{PAMS})
680 in the PAMS with several lags of SLP anomalies (0 to 3 months); (a) November (no
681 lag), (b) December, (c) January and (d) February. See Figure 1 for the description of
682 contour lines. The SST index defined in Wang et al. (2012) is shown as the grey box in
683 (a). Additional thick dashed contours represent similar lag correlations of SLPa with
684 the November SSTa index (Wang et al., 2012) for $R = 0.5$ (blue) and 0.4 (red).

685 Figure 4. Observed lead–lag correlation between the monthly meridional wind
686 anomalies in PAMS (110°E – 140°E , 5°N – 25°N) and (a) the SLP_S and (b) $-\text{PC2}$, respectively.
687 The SLP_S represents the SLP anomalies in the grey box shown in Figure 1c. See text for
688 the detailed definition of the SLP_S and $-\text{PC2}$. $p < 0.05$ is shaded.

689 Figure 5. (a) Time series of November equatorward meridional wind anomalies
690 ($-V_{\text{PAMS}}$) in the PAMS (superimposed by January $-\text{SLP}_S$ and March PC2, respectively).
691 All time series are aligned with their corresponding months. Black solid (dashed) lines
692 label the El Niño (La Niña) years in December. (b) 11-yr running correlation
693 coefficient between November V_{PAMS} and January SLP_S . The horizontal line indicates
694 the 95% confidence level.

695 Figure 6. Correlation maps between November meridional wind anomalies (V_{PAMS}) in
696 PAMS and the geopotential height anomalies (top three panels: 200 hPa, 500 hPa and
697 1000 hPa, respectively) for November and two lag months (December and January).
698 The same correlation maps with latent heat flux at different lag months are shown at the
699 bottom panel. Only $p < 0.05$ is shown.

700 Figure 7. Same as Figure 6 but for the vertical velocity (Ω) anomalies in November (no
701 lag) and January (two month lag).

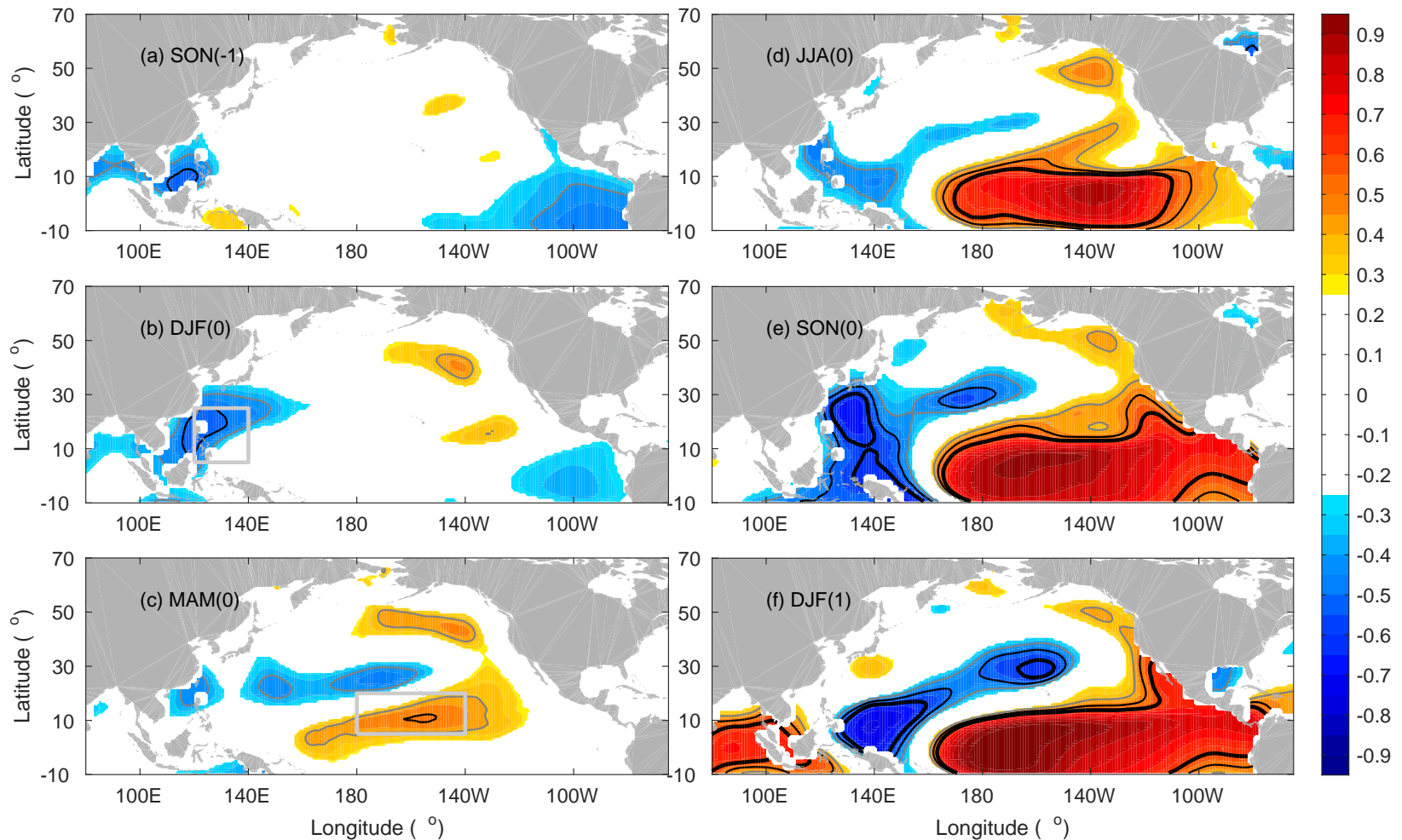
702 Figure 8: Same as Fig. 6 but for the November OLR anomalies and meridional wind
703 anomalies. Regressions of surface winds onto the V_{PAMS} are superimposed. $p < 0.05$ is
704 shown as black.

705 Fig. 9. The stationary Rossby wave ray trajectories (green curves) with initialed zonal
706 wavenumber $k = 2, 4, 6$ (upper, middle, bottom) under the NDJ (right panels) 500 hPa
707 climatological flows (gray vectors) for the period of 1965-2012. Red forks over pink
708 shadings denote wave source arrays over the PAMS region (110°E - 140°E , 5°N - 30°N).

709 Fig. 10. Same as Fig.9, but for the 250 hPa climatological flows.

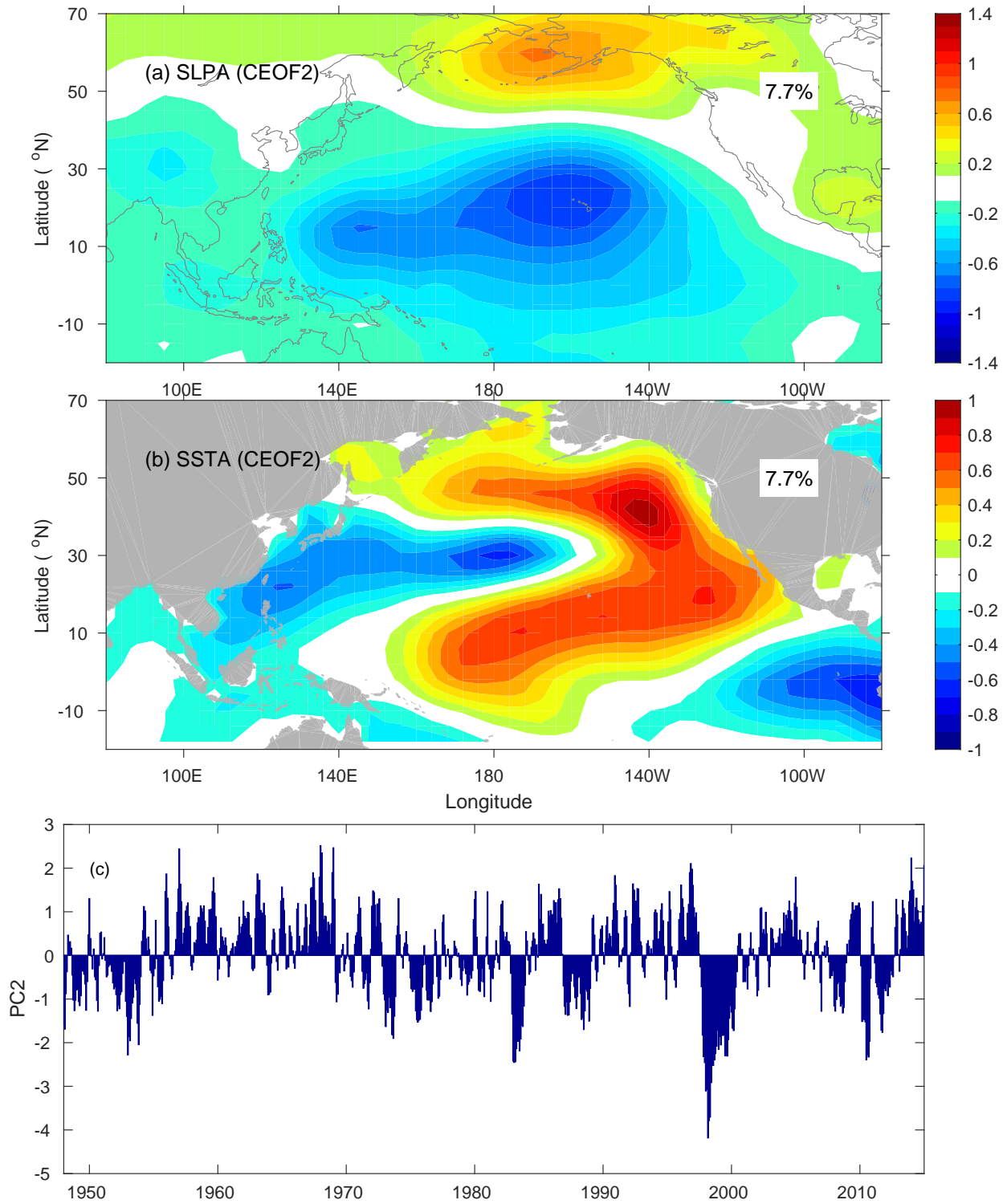
710 Fig. 11. Same as Fig.9, but for the 850 hPa climatological flows.

711 Figure 12. Schematic diagram of the dynamical processes linking the PAMS origin and
712 the onset of El Niño. The SLP (contours) and SST (color) anomalies associated with
713 CEOF2 patterns are overlaid.



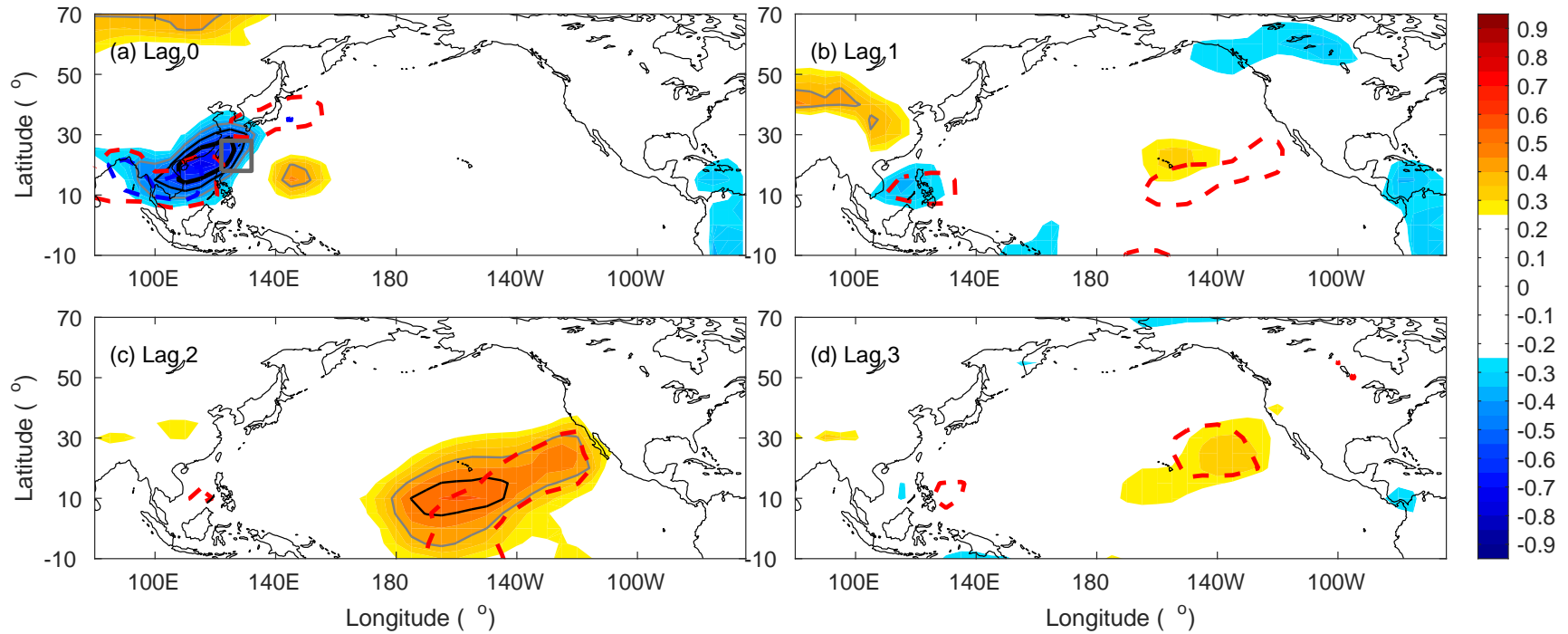
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2 Figure 1. The correlation of Pacific SST anomalies (based on 3-month running mean) and January(1) Niño4 index at different lags from 1948–
 3 2015. The Niño4 index has lags of (a) 15 months; (b) 12 months; (c) 9 months; (d) 6 months; (e) 3 months; and (f) 0 months (no lag). Only $p < 0.05$
 4 is shown. The contour lines correspond to $R=0.6$ (bold black), 0.5 (thin black) and 0.4 (grey). The grey box in (b) shows the specific PAMS region and
 5 the grey box in (c) shows the southern branch of the VM.



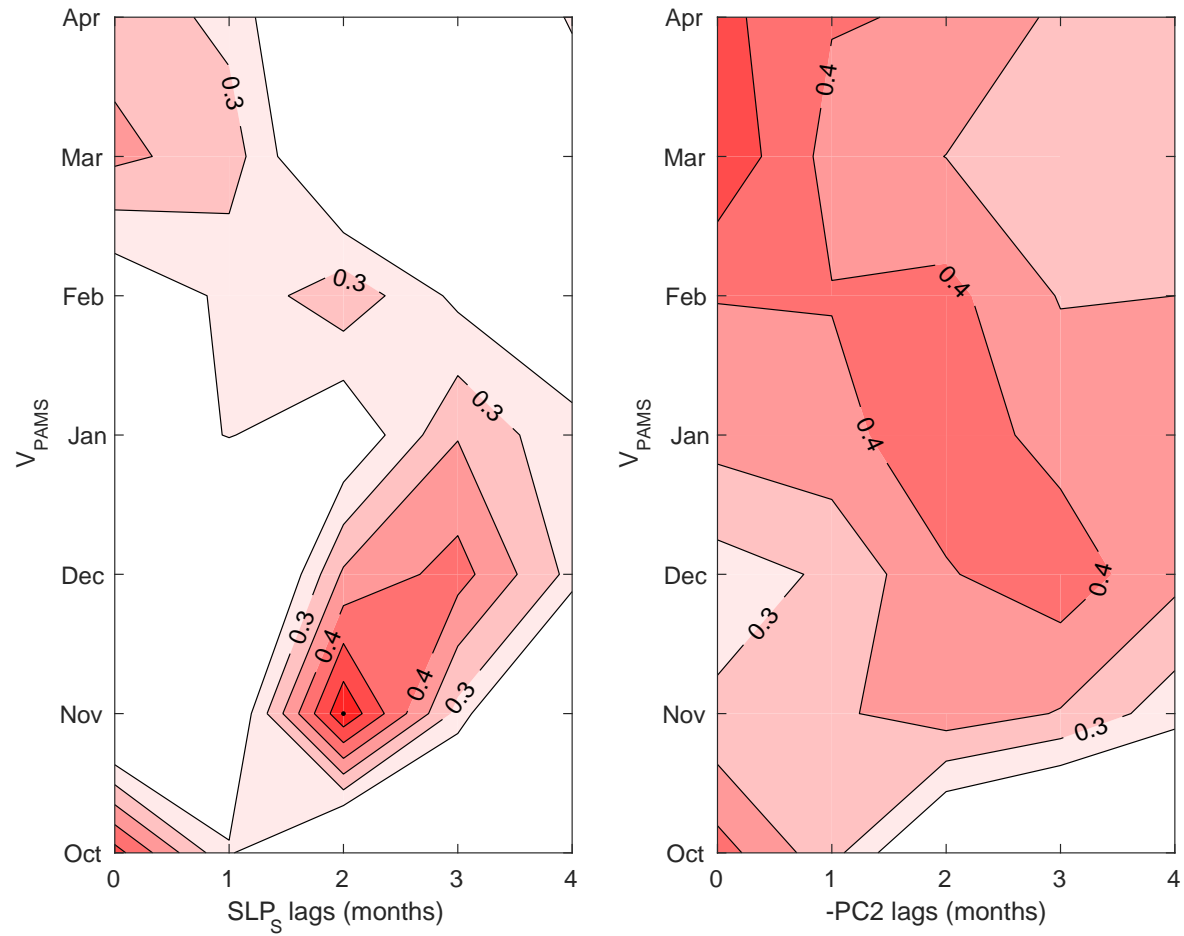
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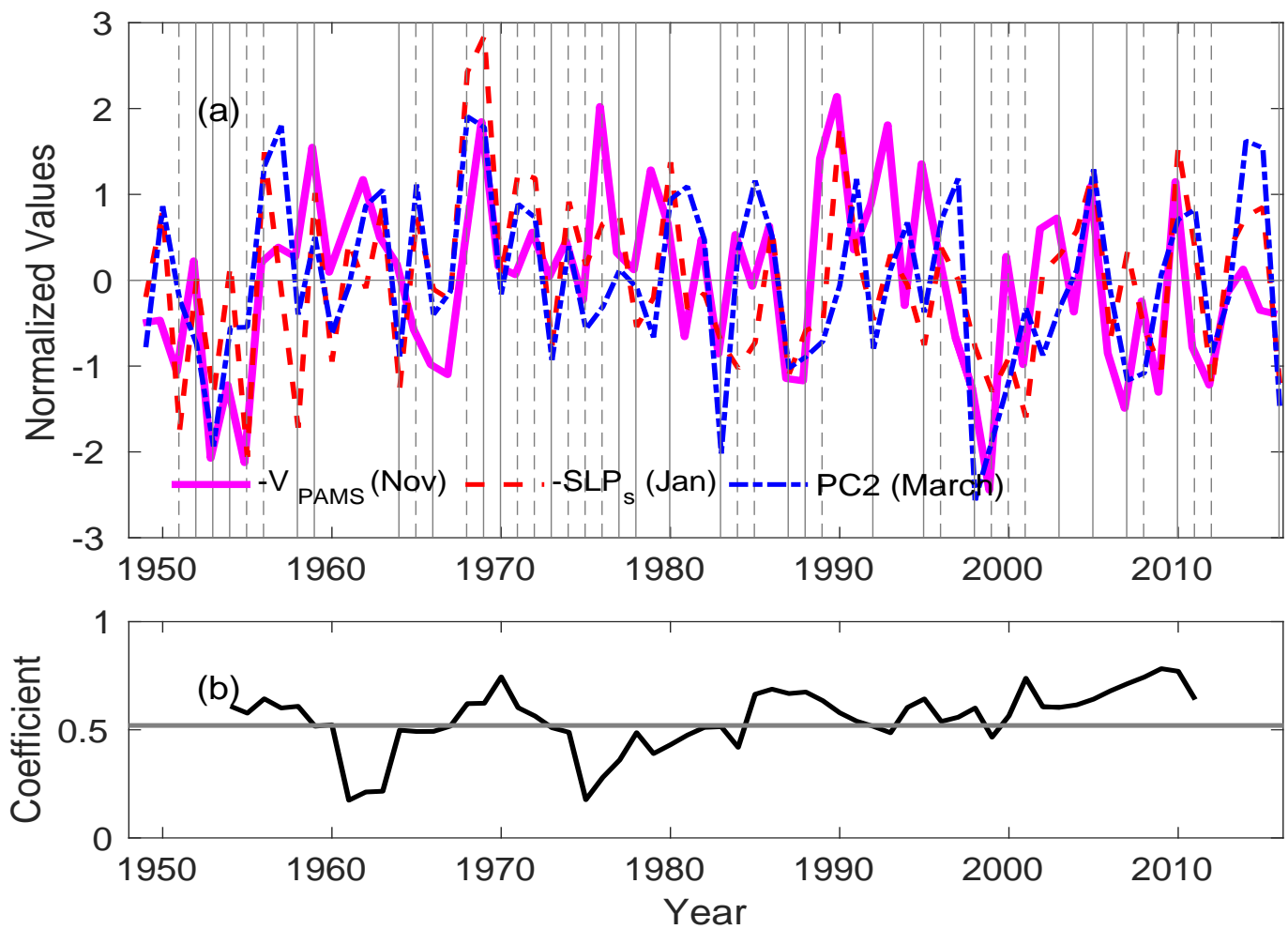
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13 Figure 3. Correlation map of the November surface meridional wind anomaly (V_{PAMS}) in the PAMS with several lags of SLP anomalies (0 to 3
14 months); (a) November (no lag), (b) December, (c) January and (d) February. See Figure 1 for the description of contour lines. The SST index
15 defined in Wang et al. (2012) is shown as the grey box in (a). Additional thick dashed contours represent similar lag correlations of SLPa with the
16 November SSTa index (Wang et al., 2012) for $R=0.5$ (blue) and 0.4 (red).



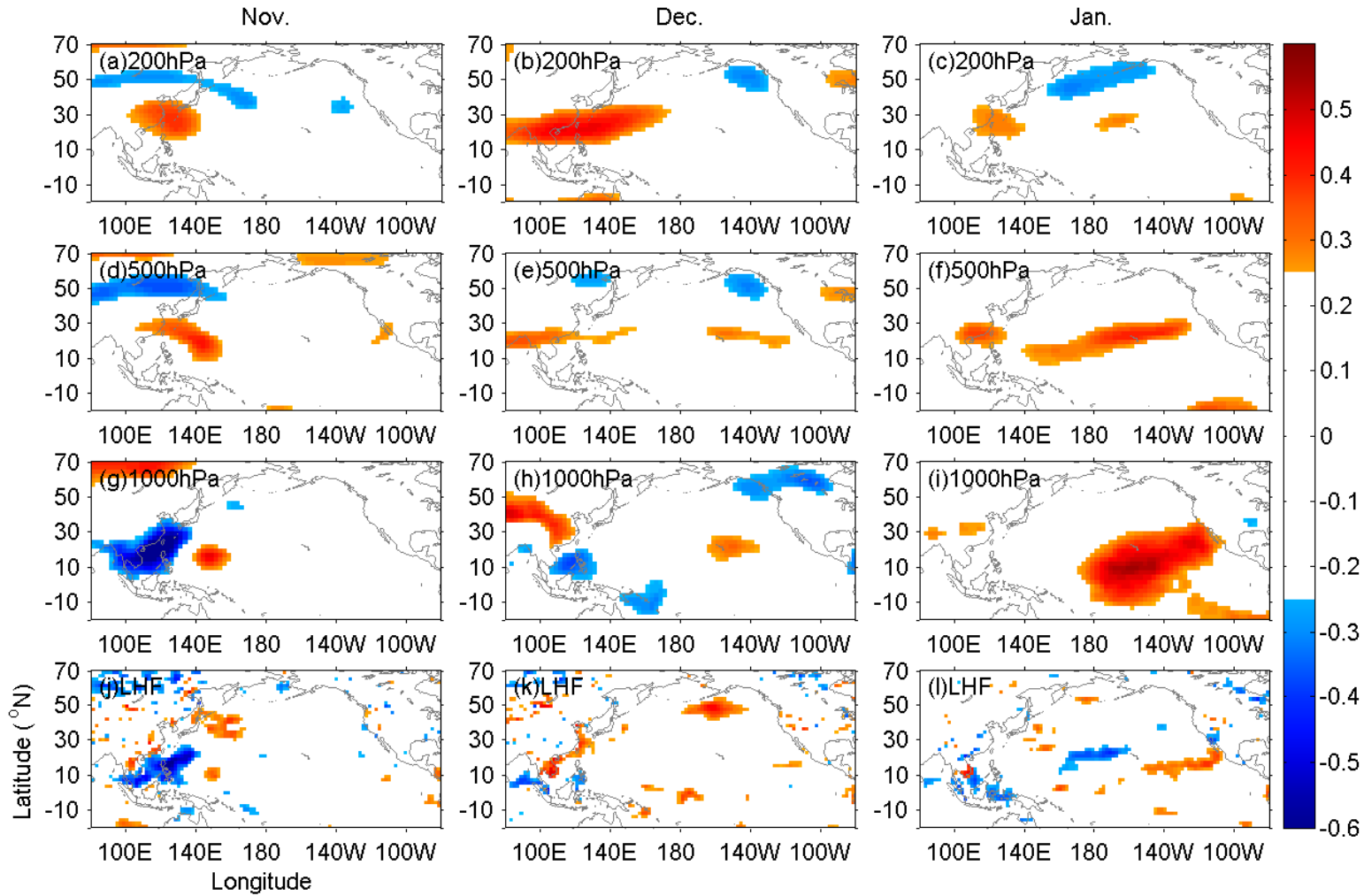
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18 Figure 4. Observed lead-lag correlation between the monthly meridional wind anomalies in PAMS (110~140°E, 5~25°N) and (a) the SLP_s and (b)
 19 $-PC2$, respectively. The SLP_s represents the SLP anomalies in the grey box shown in Figure 1c. See text for the detailed definition of the SLP_s and
 20 $-PC2$. $p < 0.05$ is shaded.



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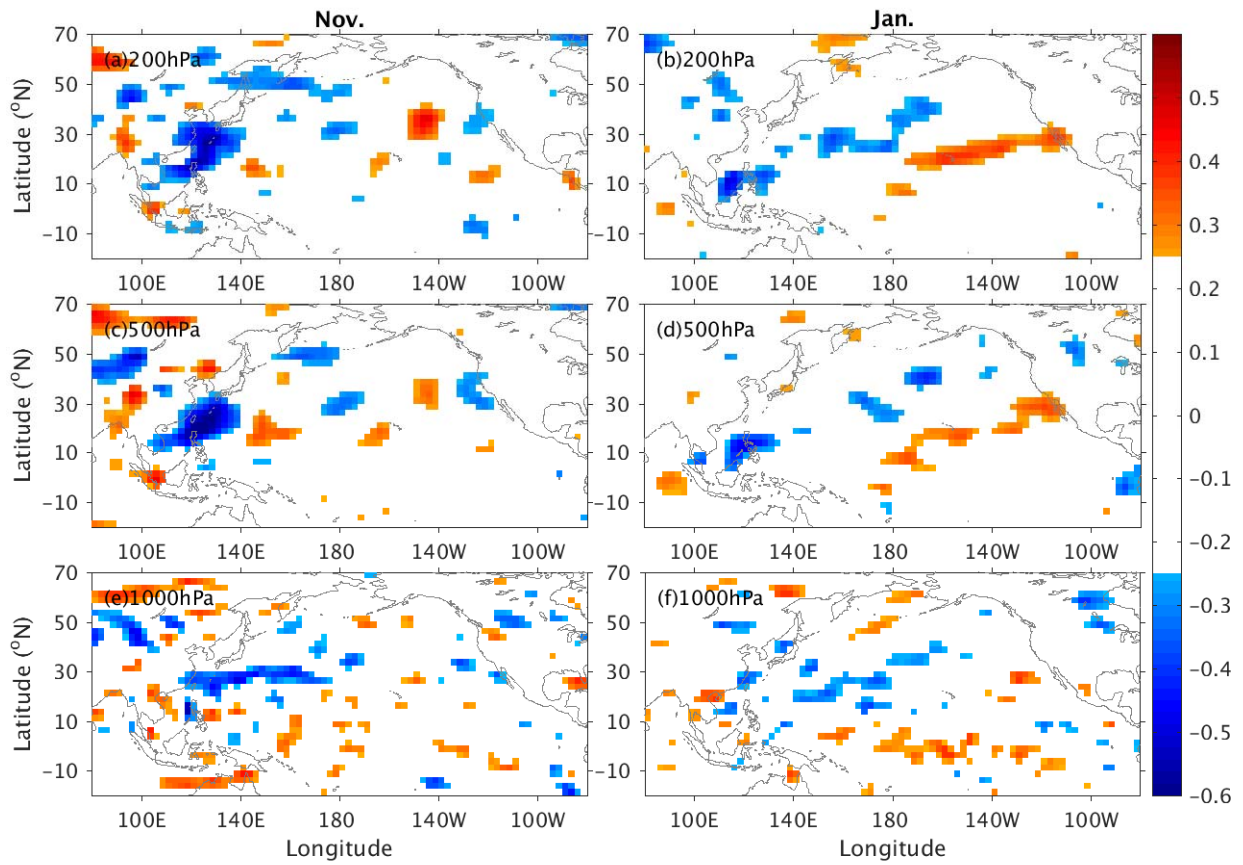
22 Figure 5. (a) Time series of November equatorward meridional wind anomalies ($-V_{PAMS}$) in the PAMS (superimposed by January $-SLP_s$ and
 23 March PC2, respectively). All time series are aligned with their corresponding months. Black solid (dashed) lines label the El Niño (La Niña)
 24 years in December. (b) 11-yr running correlation coefficient between November V_{PAMS} and January SLP_s . The horizontal line indicates the 95%
 25 confidence level.



26

27 Figure 6. Correlation maps between November meridional wind anomalies (V_{PAMS}) in PAMS and the geopotential height anomalies (top three
 28 panels: 200 hPa, 500 hPa and 1000 hPa, respectively) for November and two lag months (December and January). The same correlation maps with
 29 latent heat flux at different lag months are shown at the bottom panel. Only $p < 0.05$ is shown.

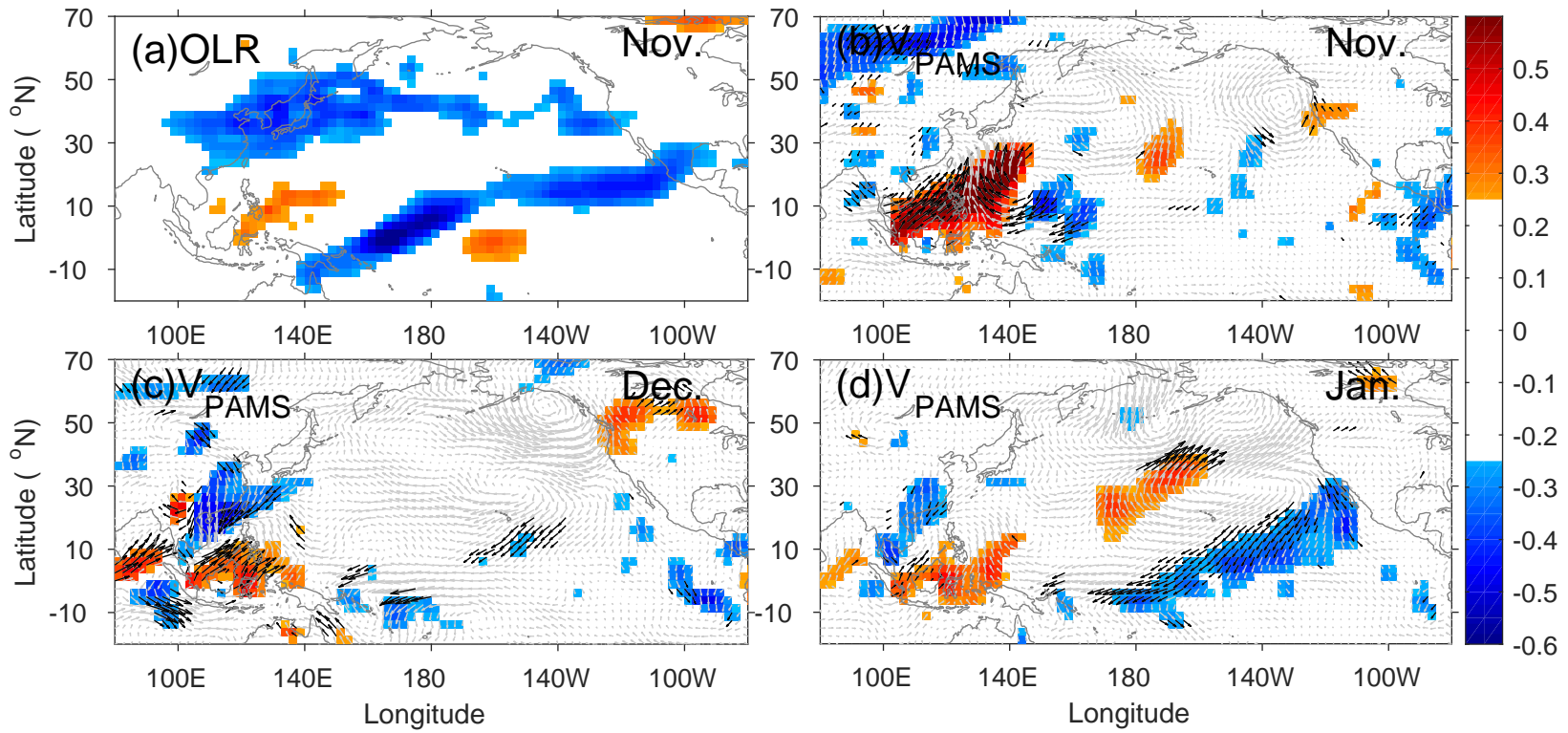
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31

32 Figure 7. Same as Figure 6 but for the vertical velocity (Ω) anomalies in November (no lag) and January
33 (two month lag).

34

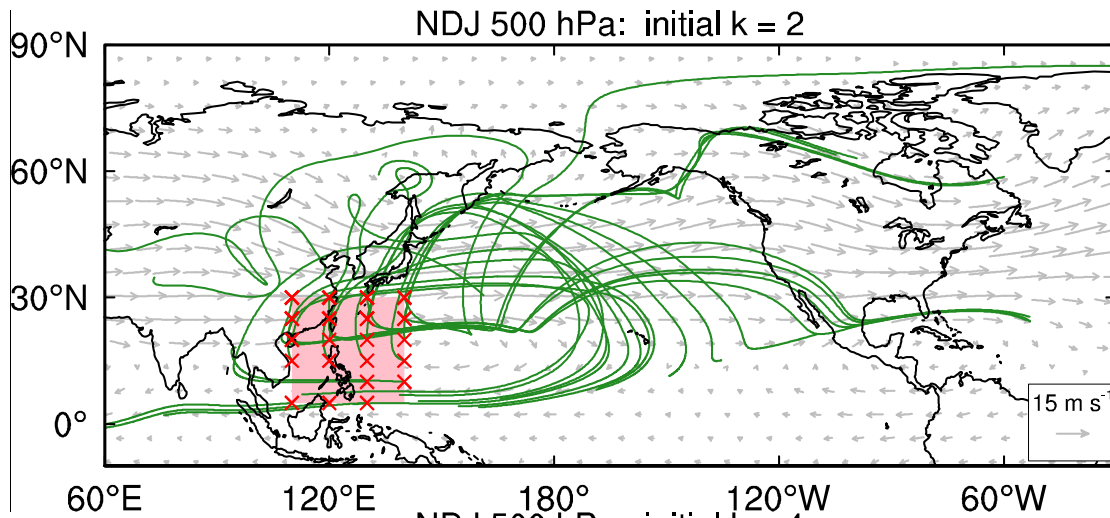


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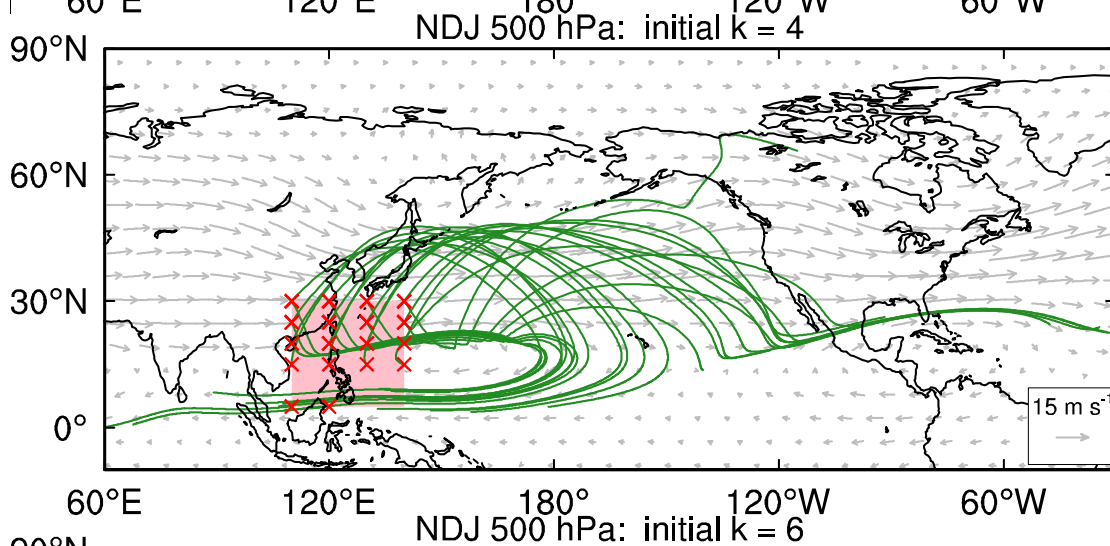
36 Figure 8: Same as Fig. 6 but for the November OLR anomalies and meridional wind anomalies. Regressions of surface winds onto the V_{PAMS} are
 37 superimposed. $p < 0.05$ is shown as black.

39

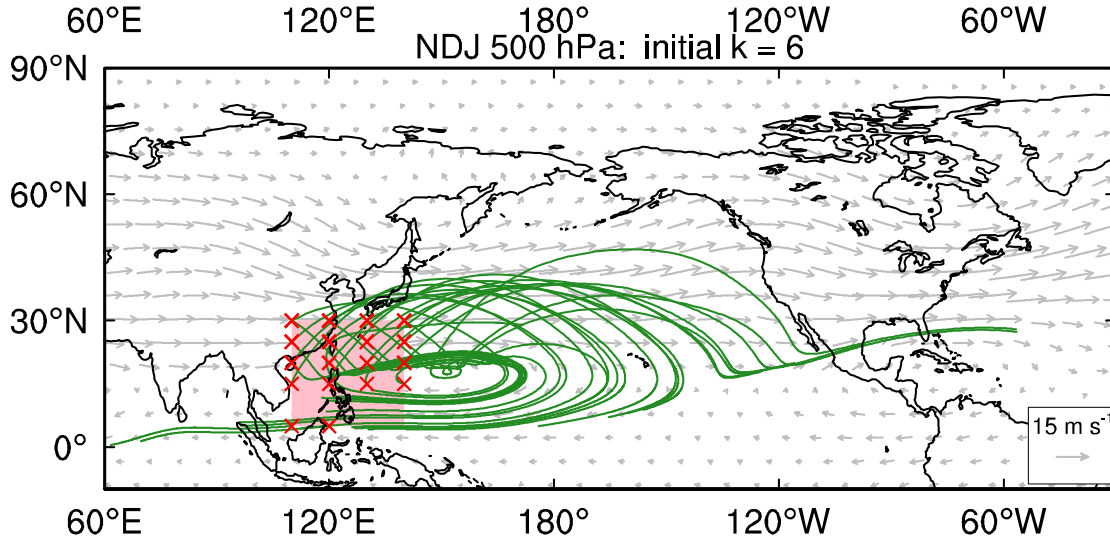
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41



42



43 Fig. 9. The stationary Rossby wave ray trajectories (green curves) with initial zonal wavenumber $k = 2$,
 44 4, 6 (upper, middle, bottom) under the NDJ (right panels) 500 hPa climatological flows (gray vectors) for
 45 the period of 1965-2012. Red forks over pink shadings denote wave source arrays over the PAMS region
 46 (110°E-140°E, 5°N-30°N).

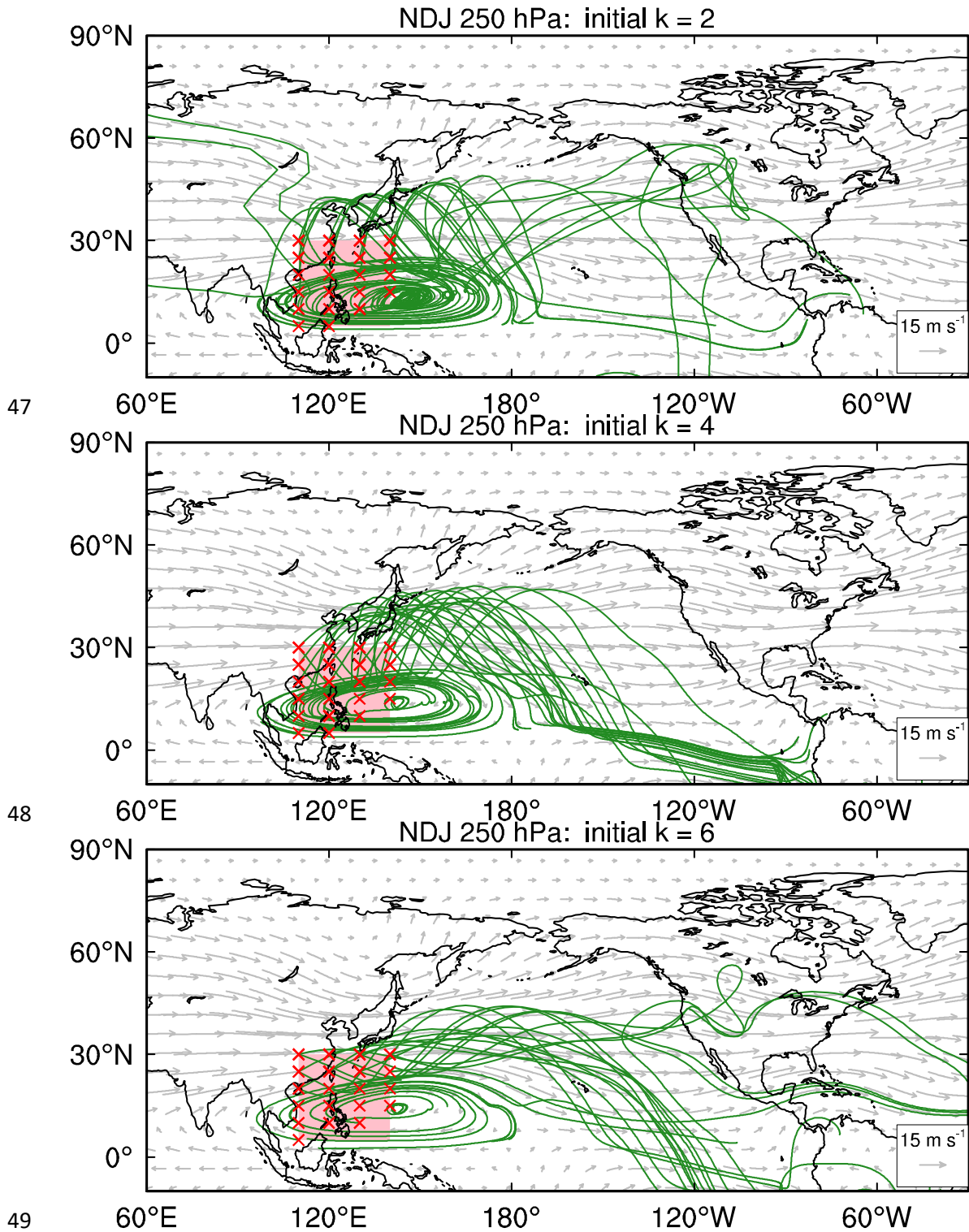
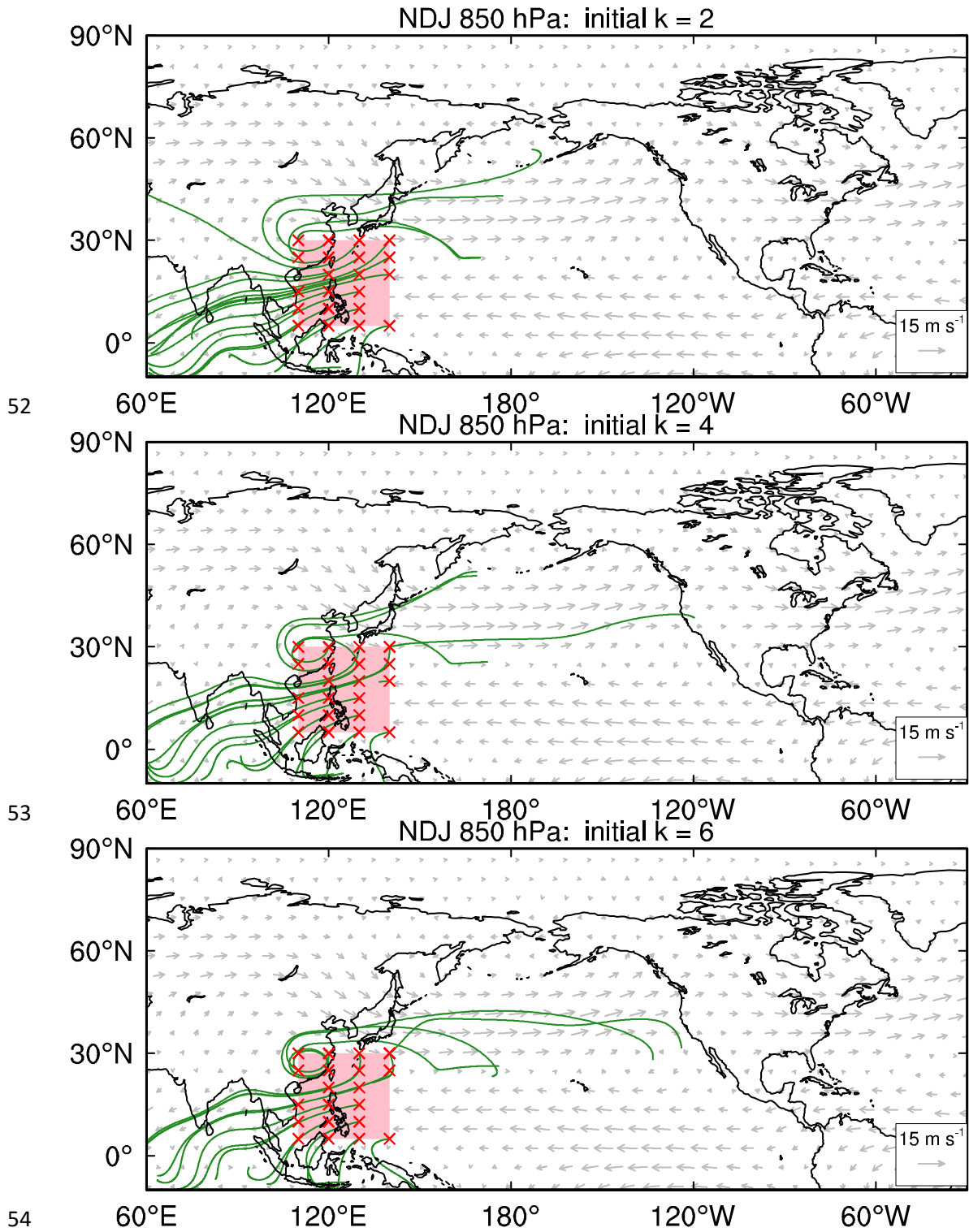


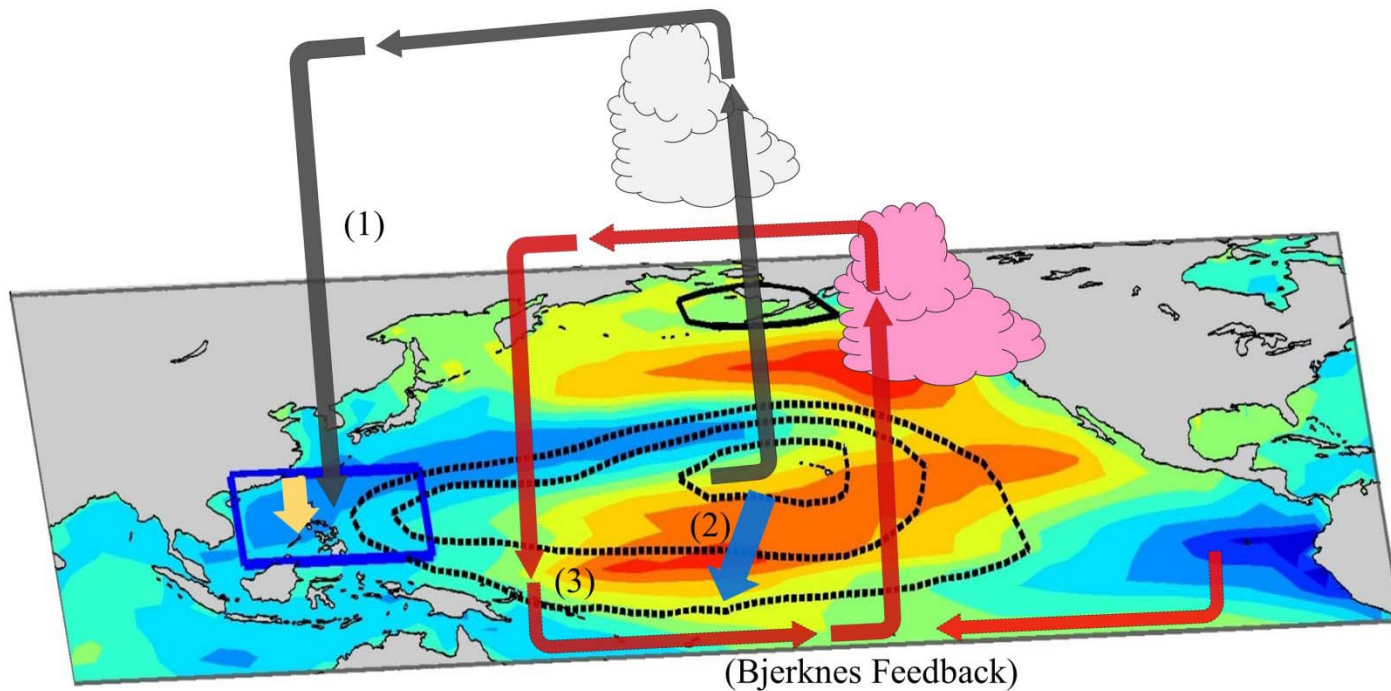
Fig. 10. Same as Fig.9, but for the 250 hPa climatological flows.

51



55 Fig. 11. Same as Fig.9, but for the 850 hPa climatological flows.

56



- (1) Win(0)-spr(0): equatorward winds in PAMS region (heat flux) → reduced deep convection in PAMS → northeastward propagation → enhanced deep convection near the south branch of the VM (south lobe of the NPO) → enhanced negative phase of NPO
- (2) Spr(0)-Sum(0): negative phase of NPO forced VM/MM (CEOF2)
- (3) Sum(0)-win(1): enhanced westerly wind in the tropical-subtropical WNP → El Niño

57
 58 Figure 12. Schematic diagram of the dynamical processes linking the PAMS origin and the onset of El Niño. The SLP (contours) and SST (color)
 59 anomalies associated with CEOF2 patterns are overlaid.